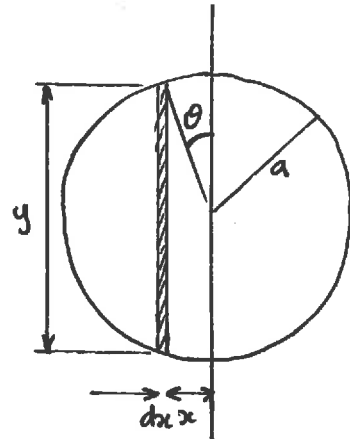


CLASSICAL MECHANICS - exercise sheet 5.

1. Part (a) is perhaps most easily deduced from the answer to part (b) using the perpendicular axis theorem but the direct solution is as follows.



(a) Consider the slice shown:

$$\text{area} = y \, dx$$

\Rightarrow if density (mass per unit area) is ρ ,

$$\text{mass of element } dm = \rho y \, dx$$

$$\text{moment of inertia } dI = \rho y \, dx \, x^2$$

and $(y/2)^2 = a^2 - x^2$

$$\Rightarrow M = \int dm = \int_{-a}^a \rho y \, dx = \int_{-a}^a \rho \sqrt{4(a^2 - x^2)} \, dx = \pi a^2 \rho \quad (\text{more directly})$$

$$I = \int dI = \int_{-a}^a \rho y x^2 \, dx = \int_{-a}^a \rho \sqrt{4(a^2 - x^2)} x^2 \, dx$$

It is more convenient to write this as a function of θ , where $x = a \sin \theta$, $y = 2a \cos \theta$, $dx = a \cos \theta \, d\theta$

$$\Rightarrow I = \int_{-\pi/2}^{\pi/2} 2\rho \sqrt{a^2 - a^2 \sin^2 \theta} a^2 \sin^2 \theta \, dx = \int_{-\pi/2}^{\pi/2} 2\rho a^4 \sin^2 \theta \cos^3 \theta \, d\theta$$

$$= \frac{\rho a^4}{2} \int_{-\pi/2}^{\pi/2} 4 \sin^2 \theta \cos^3 \theta \, d\theta = \frac{\rho a^4}{4} \int_{-\pi}^{\pi} \sin^2 2\theta \, d(2\theta) = \frac{\rho a^4}{4} \frac{2\pi}{2} \quad \text{since } \sin^2 \text{ averages } \frac{1}{2} \text{ over a cycle}$$

$$\Rightarrow \underline{I = \frac{\pi \rho a^4}{4} = \frac{1}{4} M a^2.}$$

(b) Since by symmetry $I_x = I_y$, the perpendicular axis theorem gives directly

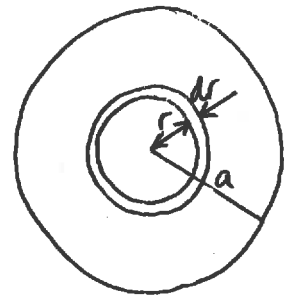
$$\underline{I_z = \frac{1}{4} M a^2 + \frac{1}{4} M a^2 = \frac{1}{2} M a^2.}$$

This may be derived directly by considering concentric rings

$$dm = 2\pi r \, dr \, \rho \quad \dots \text{hence mass per above}$$

$$dI = 2\pi r \, dr \, \rho \, r^2$$

$$\Rightarrow I = \int_0^a 2\pi \rho r^3 \, dr = 2\pi \rho \left[\frac{r^4}{4} \right]_0^a = \frac{\pi \rho a^4}{2} = \underline{\underline{\frac{1}{2} M a^2.}}$$



Given this result and the symmetry that $I_x = I_y$, it follows that $I_x = I_y = \frac{1}{2} \left(\frac{1}{2} M a^2 \right) = \underline{\underline{\frac{1}{4} M a^2.}}$

CLASSICAL MECHANICS - exercise sheet 5 - cont'd.

$$\begin{aligned} 2. a) \underline{L} &= \underline{r} \times \underline{p} = \underline{r} \times m \underline{\dot{r}} \\ &= m \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x_0 + at^2 & bt & ct^3 \\ 2at & b & 3ct^2 \end{vmatrix} \\ &= m \begin{pmatrix} 3bct^3 - bct^3 \\ 2act^4 - 3ct^2(x_0 + at^2) \\ b(x_0 + at^2) - 2abt^2 \end{pmatrix} = m \begin{pmatrix} 2bct^3 \\ -c(at^4 + 3x_0t^2) \\ b(x_0 - at^2) \end{pmatrix} \\ &\equiv m \{ 2bct^3 \hat{i} - c(at^4 + 3x_0t^2) \hat{j} + b(x_0 - at^2) \hat{k} \}. \end{aligned}$$

$$b) \underline{F} = m \underline{\ddot{r}} = \underline{m} \{ 2a \hat{i} + 6ct \hat{k} \}.$$

$$\begin{aligned} \frac{d\underline{L}}{dt} &= m \{ 6bct^2 \hat{i} - c(4at^3 + 6x_0t) \hat{j} + b(-2at) \hat{k} \} \\ \underline{r} \times \underline{F} &= m \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x_0 + at^2 & bt & ct^3 \\ 2a & 0 & 6ct \end{vmatrix} = m \begin{pmatrix} 6bct^2 \\ 2act^3 - 6ct(x_0 + at^2) \\ -2abt \end{pmatrix} \\ &= m \begin{pmatrix} 6bct^2 \\ -c(4at^3 + 6x_0t) \\ -2abt \end{pmatrix} \end{aligned}$$

$$\equiv m \{ 6bct^2 \hat{i} - c(4at^3 + 6x_0t) \hat{j} - 2abt \hat{k} \}$$

$$\Rightarrow \underline{\frac{dL}{dt}} = \underline{r} \times \underline{F}.$$

CLASSICAL MECHANICS - exercise sheet 5 - cont'd.

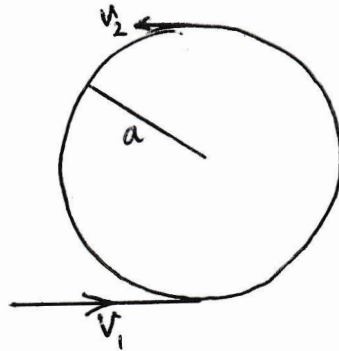
3. Energy is conserved (since none is lost)

$$\Rightarrow \frac{1}{2} m v_1^2 = \frac{1}{2} m v_2^2 + m g 2a$$

where m is the glider's mass.

Rearranging and cancelling m , gives

$$v_1^2 = v_2^2 + 4ga$$



To be "weightless" at the top of the loop, the centripetal acceleration v_2^2/a should be provided by gravity, i.e.

$$v_2^2/a = g.$$

Combining our expressions,

$$v_1^2 = ag + 4ga = 5ag$$

Hence the acceleration at the beginning of the loop will be

$$\frac{v_1^2}{a} = \underline{\underline{5g}}.$$

The wings must produce a lift force to provide this, as well as supporting the glider's weight, so the lift required will be $5mg + mg = \underline{\underline{6 \times \text{glider's weight (mg)}}}$.

CLASSICAL MECHANICS - exercise sheet 5 - cont'd.

H: a) Consider a ring of radius r , cross-section d and dr

$$\Rightarrow \text{mass } dm = 2\pi r d dr \rho$$

distance from measurement point P is

$$s = \sqrt{h^2 + r^2}$$

The gravitational forces from each element of the ring will have the same vertical component but horizontal components that cancel, so the total will be vertically downward with magnitude

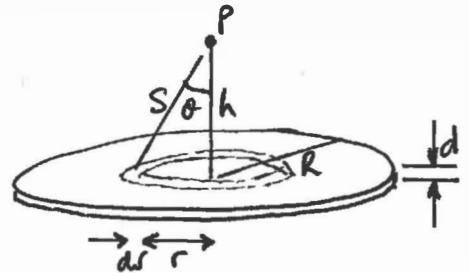
$$\begin{aligned} dF &= \frac{Gm dm \cos\theta}{s^2} = \frac{Gm 2\pi d \rho r dr}{h^2 + r^2} \frac{h}{s} \\ &= Gm 2\pi d \rho h \frac{r dr}{(h^2 + r^2)^{3/2}} \end{aligned}$$

\Rightarrow total force for the disc will be

$$\begin{aligned} F &= \int_{r=0}^{r=R} dF = Gm 2\pi d \rho h \int_0^R \frac{r}{(h^2 + r^2)^{3/2}} dr \\ &= Gm 2\pi d \rho h \left[-(h^2 + r^2)^{-1/2} \right]_0^R \\ &= Gm 2\pi d \rho h \left(\frac{1}{h} - \frac{1}{\sqrt{h^2 + R^2}} \right) \\ &= 2\pi G m \rho d \left(1 - \frac{h}{\sqrt{h^2 + R^2}} \right) \end{aligned}$$

hence the gravitational field F/m will be

$$\underline{\underline{g_D = 2\pi G \rho d \left(1 - \frac{h}{\sqrt{h^2 + R^2}} \right)}}.$$



CLASSICAL MECHANICS - exercise sheet 5 - cont'd.

4.b) If the room is sufficiently large, we may assume the result for $R \gg h$, which gives $g_D = 2\pi C \rho d$

$$\text{With the values given } g_D = 2\pi \cdot 6.67 \times 10^{-11} \text{ Nm}^2 \text{kg}^{-2} \cdot 1025 \text{ kg m}^{-3} \cdot 0.15 \text{ m} \\ = 64 \times 10^{-8} \text{ ms}^{-2}.$$

The pendulum period is given by $T = 2\pi \sqrt{\frac{l}{g}}$ for pendulum length l

$$\Rightarrow \frac{dT}{dg} = 2\pi l^{\frac{1}{2}} \left(-\frac{1}{2} g^{-\frac{3}{2}}\right)$$

$$= \frac{-1}{2g} T$$

$$\Rightarrow \frac{dT}{T} = \frac{-1}{2} \frac{dg}{g} = \frac{-1}{2} \frac{64 \times 10^{-8} \text{ ms}^{-2}}{9.8 \text{ ms}^{-2}} \quad \text{with the values given}$$

In one year, the clock will therefore gain (since the period is reduced)

$$\Delta T = \frac{1}{2} \frac{64 \times 10^{-8}}{9.8} (365 \times 24 \times 3600 \text{ s}) \\ = \underline{\underline{0.1 \text{ s}}}$$

c) If the seam of iron ore is wide, we may again assume $R \gg h$ and obtain, as above,

$$\Delta T = \frac{1}{2} \frac{(3000-2700) 2\pi \cdot 6.67 \times 10^{-11} \cdot 120}{9.8} (365 \times 24 \times 3600) \\ = \underline{\underline{24 \text{ s/year}}}$$

Since the iron ore is denser than the typical Earth's crust, the clock will again run faster, with a shorter oscillation period, and gain 24 s/year.