

Wave Physics

PHYS 2023

Tim Freegarde



Wave propagation

waves are

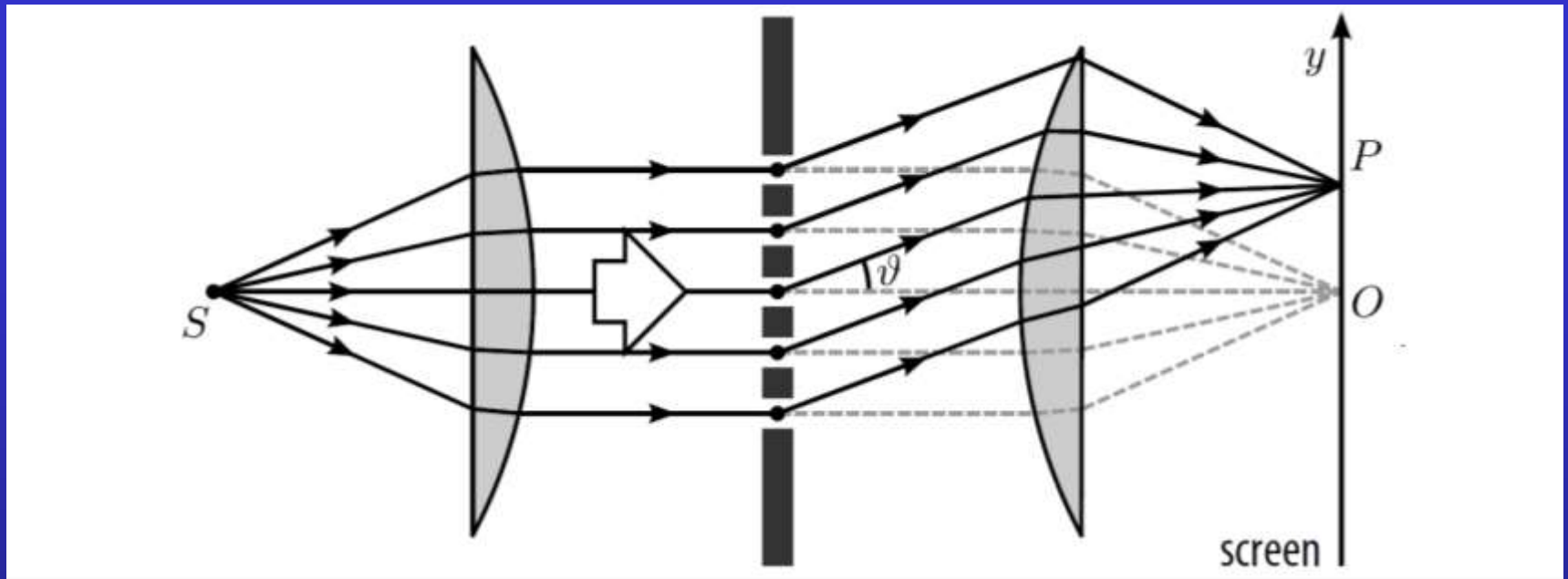
- collective bulk disturbances, in which
- motion is a delayed response
- to neighbouring motions

when propagation follows multiple routes

- the amplitudes are added
- waves propagate via all possible routes

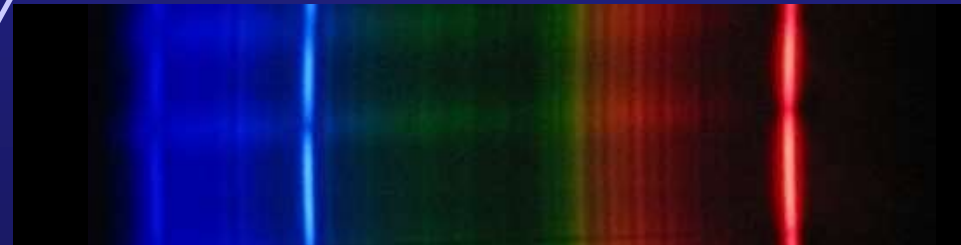
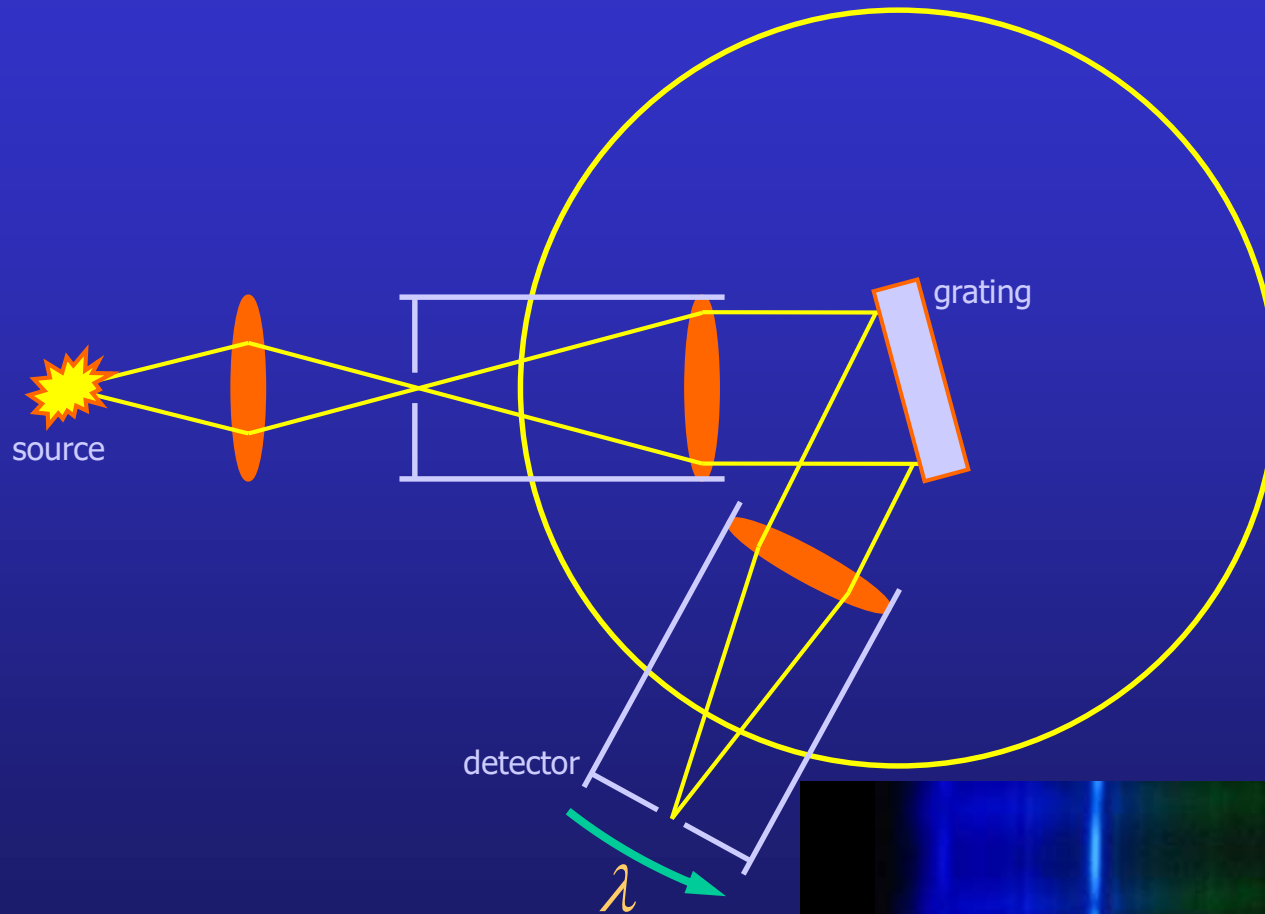


Fraunhofer diffraction



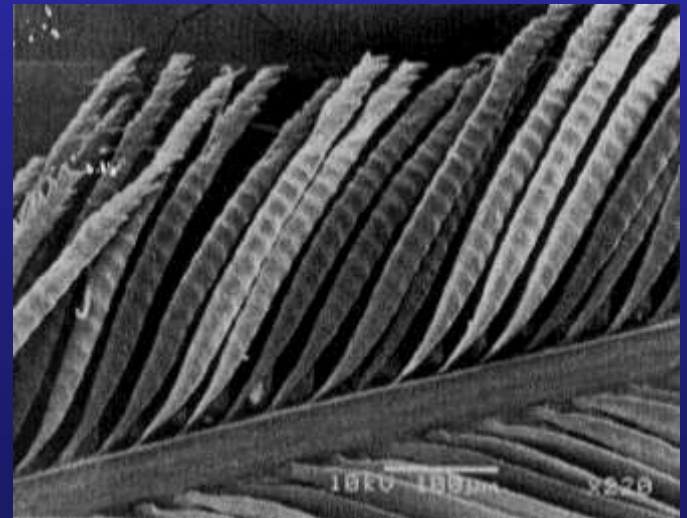
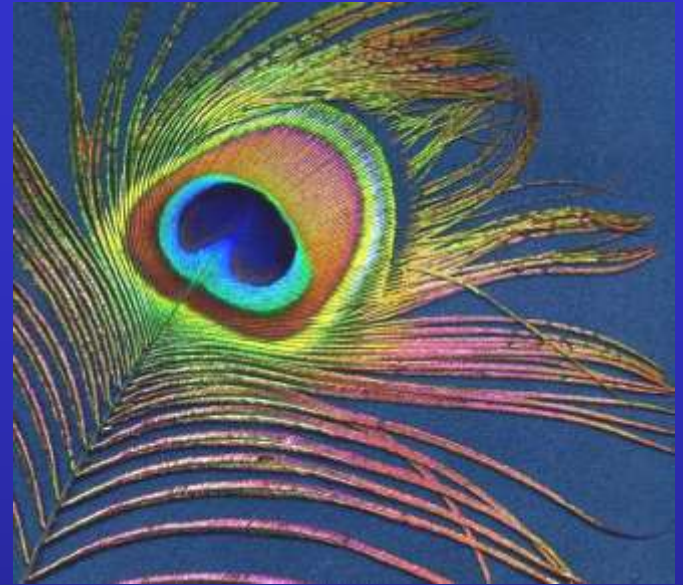
- observed at **infinity** with **plane-wave** illumination
- observed using **imaging lens** and **plane-wave** illumination
- observed in **image plane** of source
- path length phase is **linear function** of **transverse coordinate** in object

Grating spectrometer



Diffraction

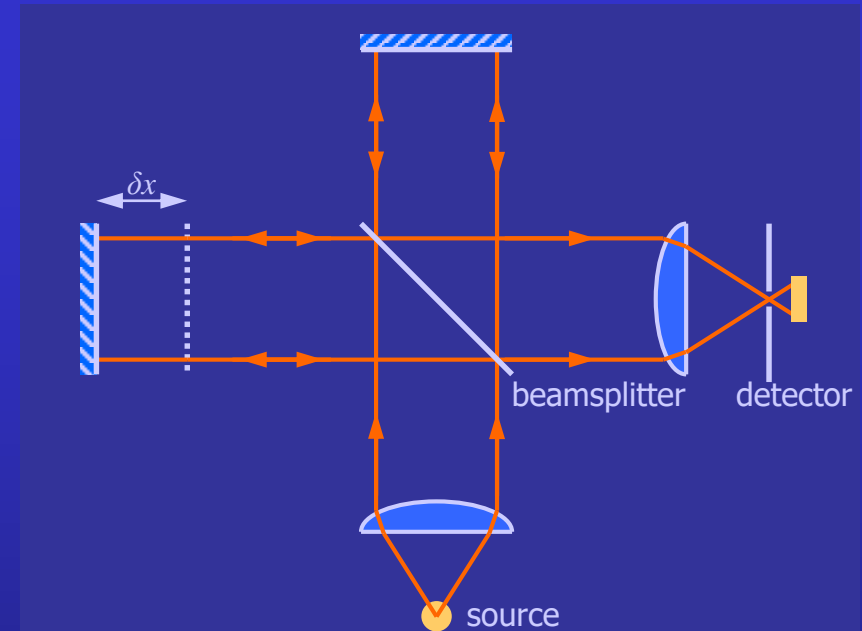
- iridescence of feathers (Grimaldi, 1665)
- interference by **division of wavefront**



S Yoshioka & S Kinoshita, *Forma* **17** 169 (2002)

Michelson interferometer

- interference by **division of amplitude**



Consider a wave propagating from the source that, after the collimating lens, is

$$\underline{E}(x, t) = \underline{E}_0 \cos(kx - \omega t)$$

This is split and sent by two paths a and b so that two components approach the focussing lens:

$$\underline{E}_a(x, t) = r t \underline{E}_0 \cos[k(x_0 + 2x_a) - \omega t] = r t \underline{E}_0 \cos[k(x_0 + x_a + x_b + (x_a - x_b)) - \omega t]$$

$$\underline{E}_b(x, t) = t r \underline{E}_0 \cos[k(x_0 + 2x_b) - \omega t] = t r \underline{E}_0 \cos[k(x_0 + x_a + x_b - (x_a - x_b)) - \omega t]$$

By using the identity $\cos(A+B) = \cos A \cos B - \sin A \sin B$, we write the superposition as

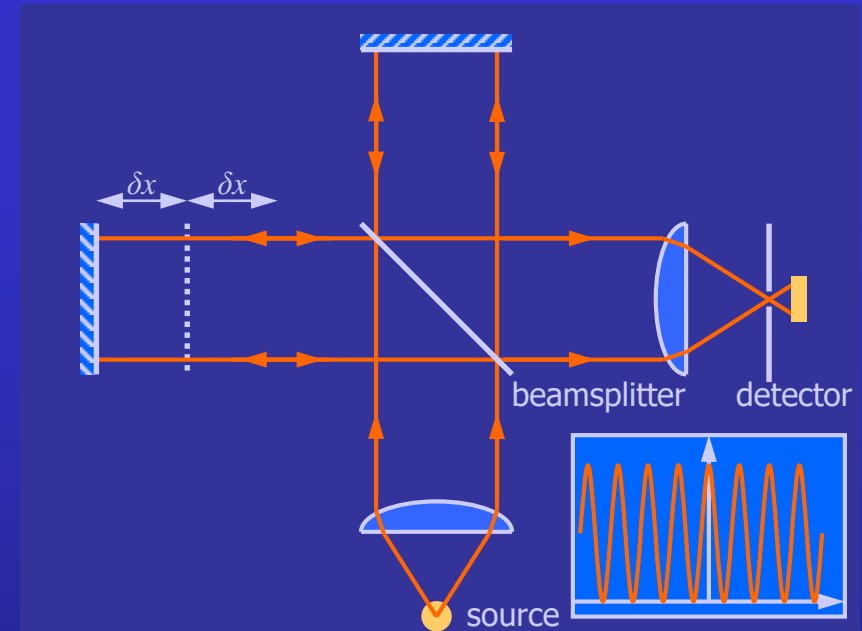
$$\underline{E}_a(x, t) + \underline{E}_b(x, t) = \underbrace{2 r t \underline{E}_0 \cos[k(x_0 + x_a + x_b) - \omega t]}_{\text{SINGLE BEAM ALONG MEAN PATH}} \underbrace{\cos[k(x_a - x_b)]}_{\text{MODULATION}}$$

The (time-averaged) intensity will be

$$I(\delta x) \propto 4 R T \cos^2[k \delta x] \quad \text{where } \delta x \equiv x_a - x_b, R \equiv r^2, T \equiv t^2$$

Michelson interferometer

- interference by **division of amplitude**



Consider a wave propagating from the source that, after the collimating lens, is

$$\underline{E}(x, t) = \underline{E}_0 \cos(kx - \omega t)$$

This is split and sent by two paths a and b so that two components approach the focussing lens:

$$\underline{E}_a(x, t) = r t \underline{E}_0 \cos[k(x_0 + 2x_a) - \omega t] = r t \underline{E}_0 \cos[k(x_0 + x_a + x_b + (x_a - x_b)) - \omega t]$$

$$\underline{E}_b(x, t) = t r \underline{E}_0 \cos[k(x_0 + 2x_b) - \omega t] = t r \underline{E}_0 \cos[k(x_0 + x_a + x_b - (x_a - x_b)) - \omega t]$$

By using the identity $\cos(A+B) = \cos A \cos B - \sin A \sin B$, we write the superposition as

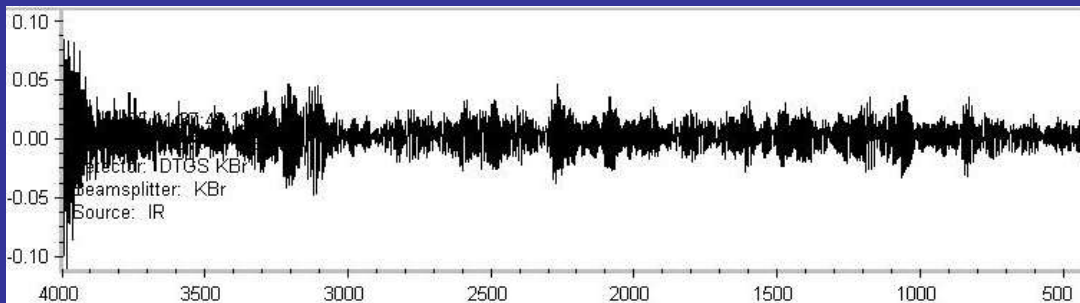
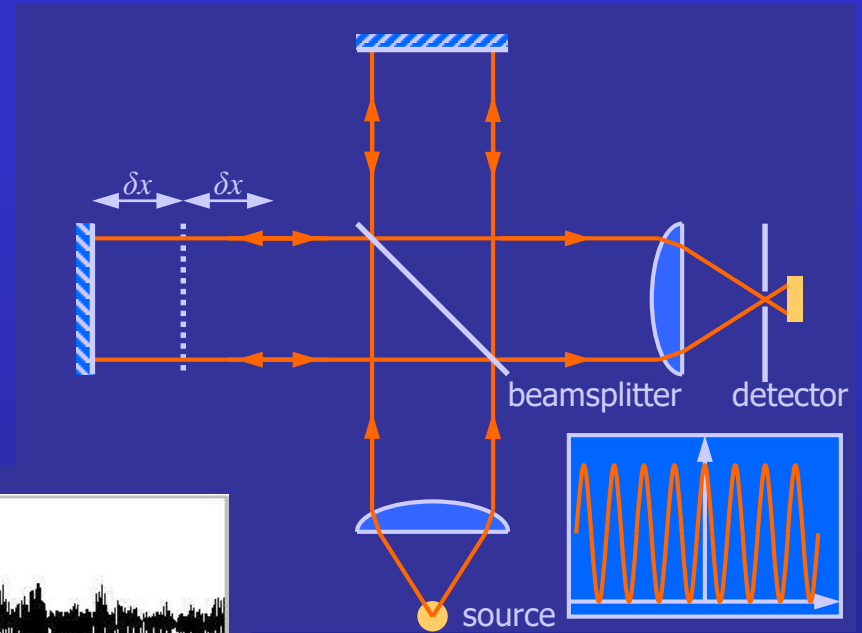
$$\underline{E}_a(x, t) + \underline{E}_b(x, t) = \underbrace{2 r t \underline{E}_0 \cos[k(x_0 + x_a + x_b) - \omega t]}_{\text{SINGLE BEAM ALONG MEAN PATH}} \underbrace{\cos[k(x_a - x_b)]}_{\text{MODULATION}}$$

The (time-averaged) intensity will be

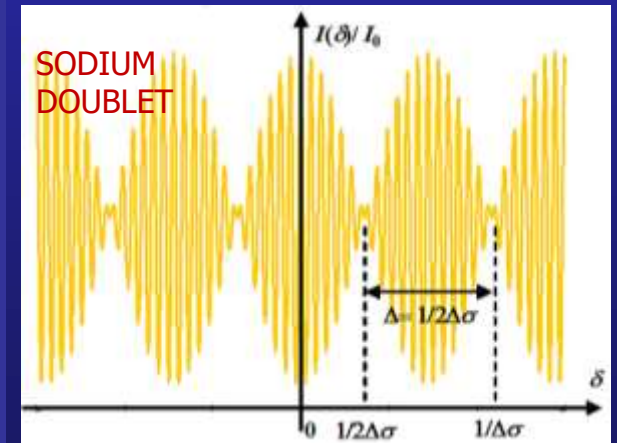
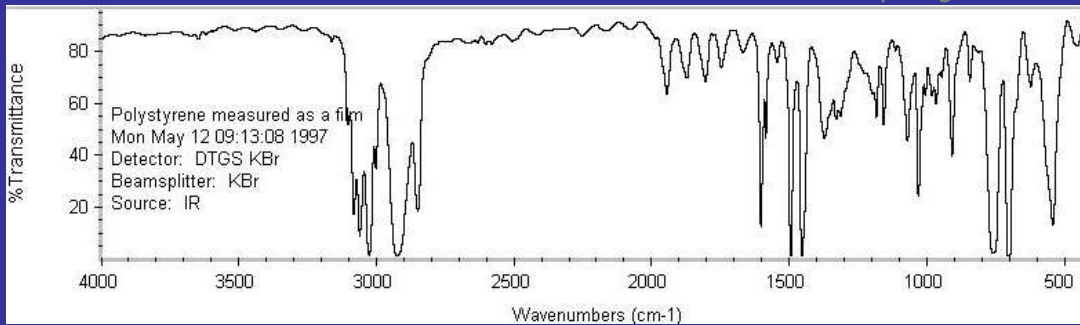
$$I(\delta x) \propto 4 R T \cos^2[k \delta x] \quad \text{where } \delta x \equiv x_a - x_b, R \equiv r^2, T \equiv t^2$$

Michelson interferometer

- interference by **division of amplitude**
- measurement of **wavelengths, spectra and lineshapes**
- FTIR: **Fourier transform infrared**



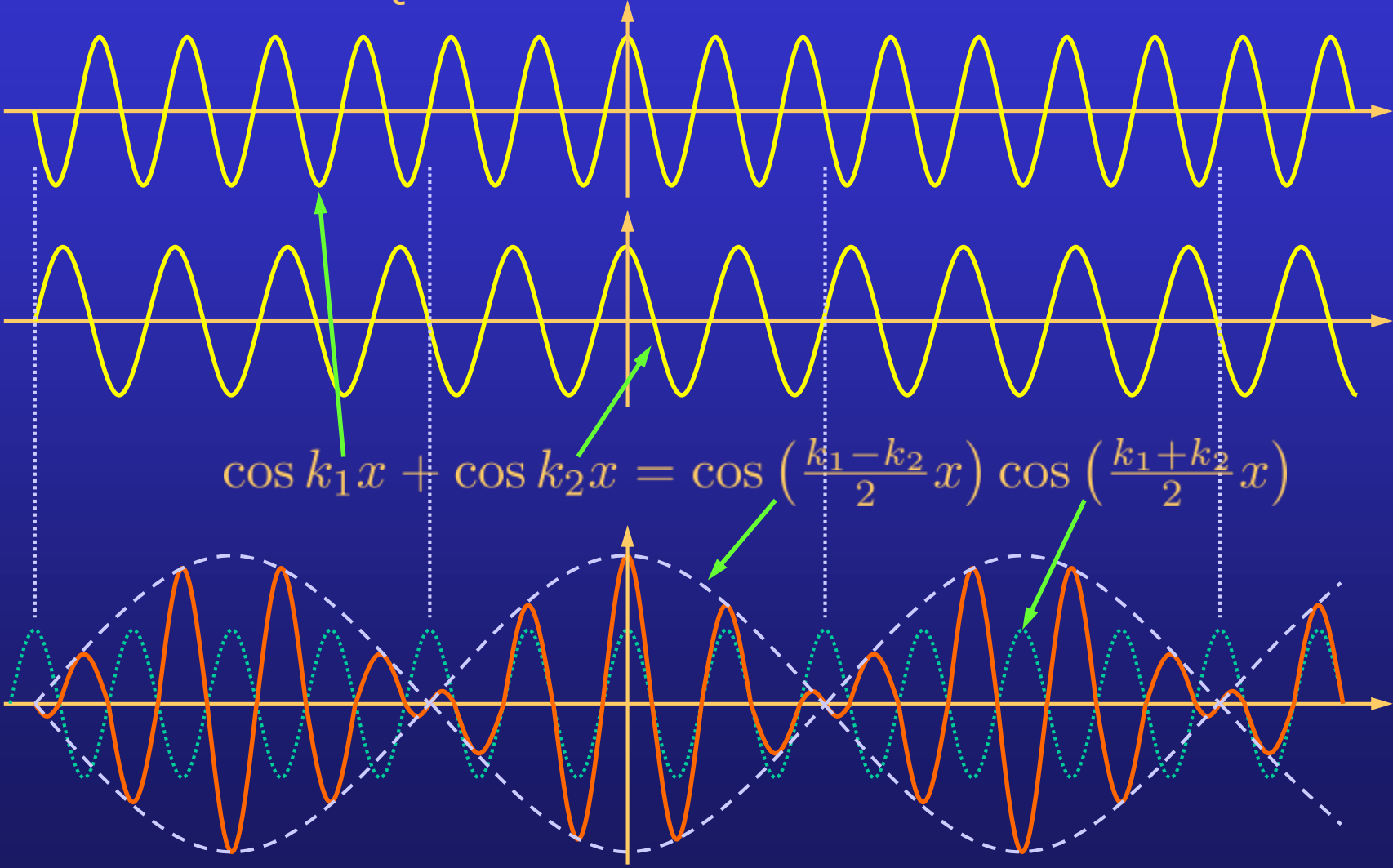
chemistry.oregonstate.edu



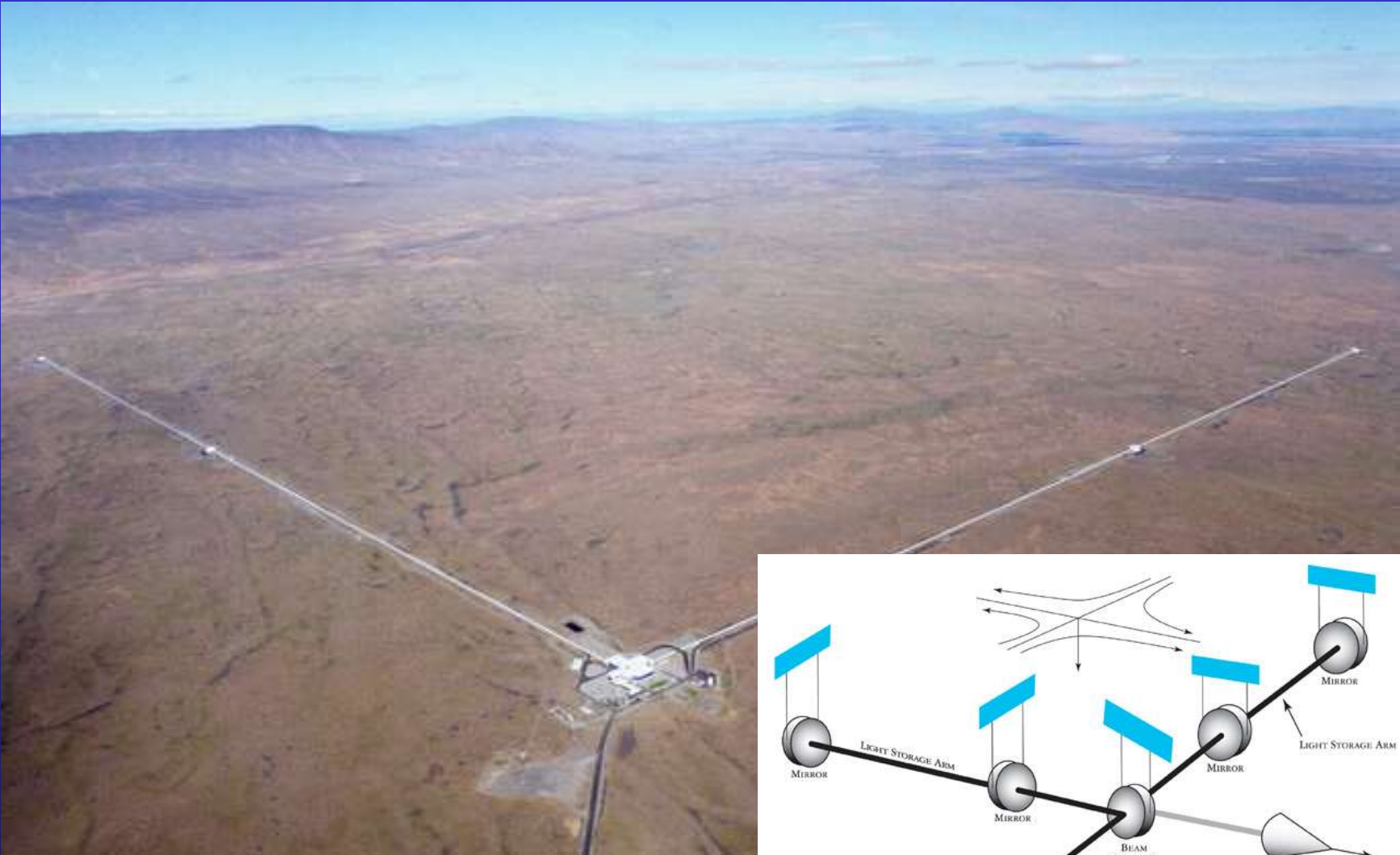
optique-ingenieur.org

Beating

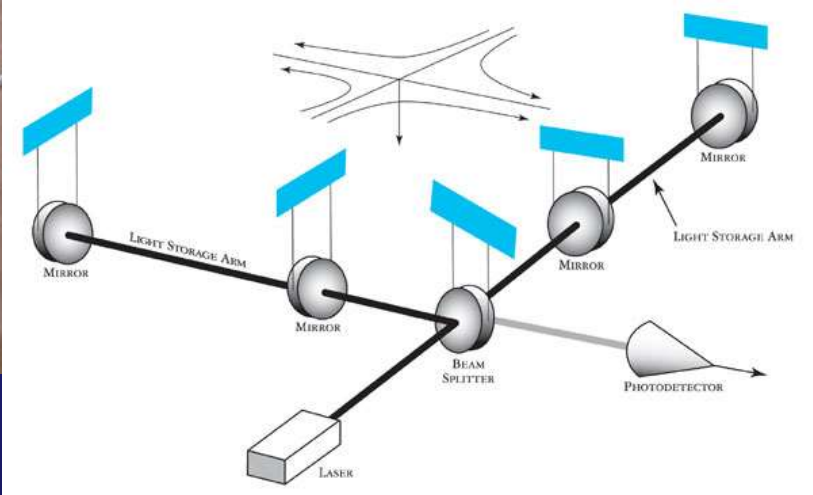
TWO DIFFERENT FREQUENCIES



LIGO

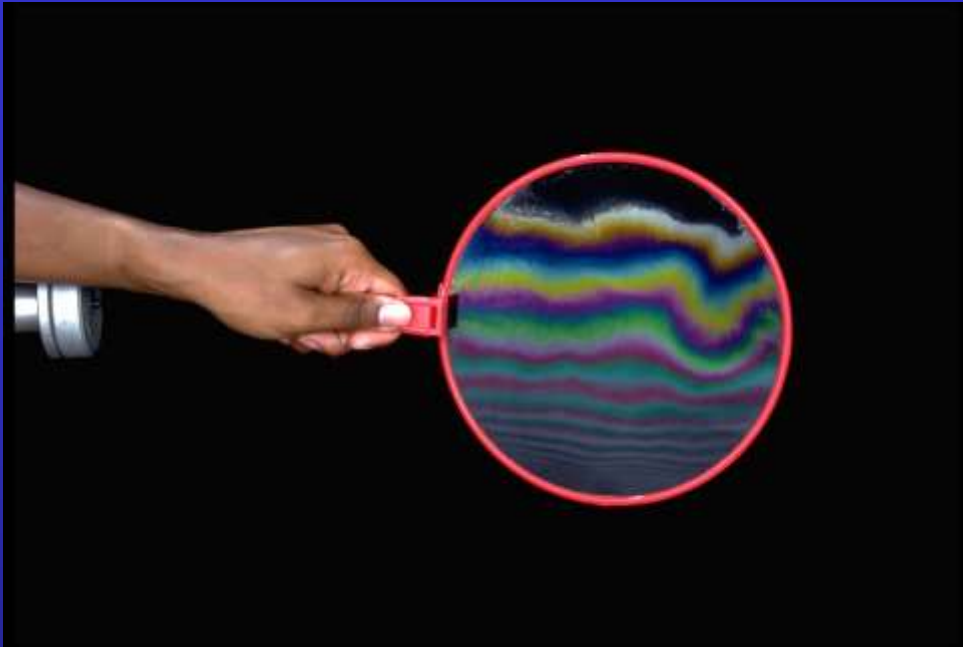


Caltech/MIT/LIGO Lab

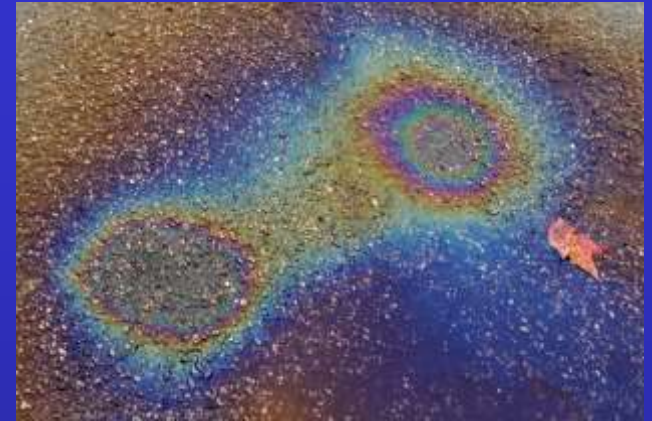


Optics & Photonics

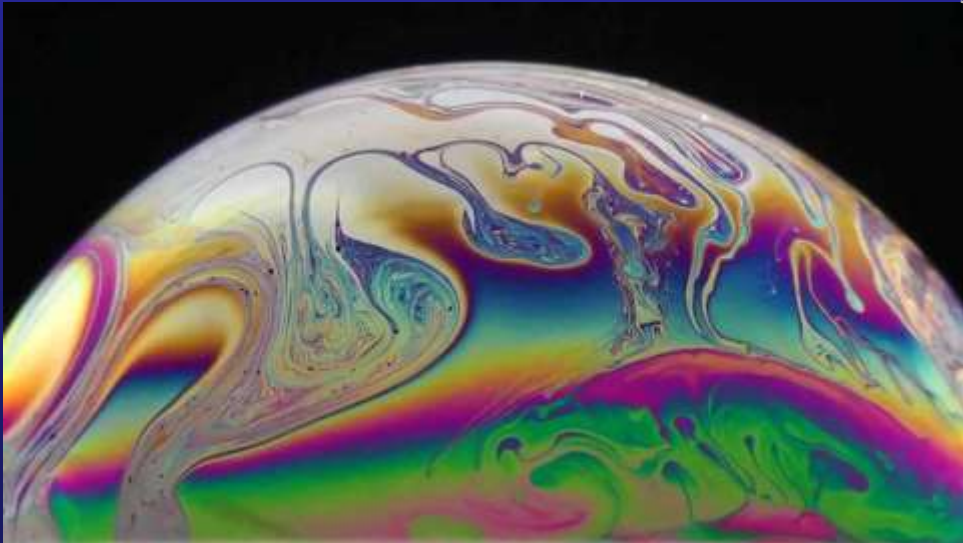
Thin films



Susan Schwartzberg



Caroline Mockett

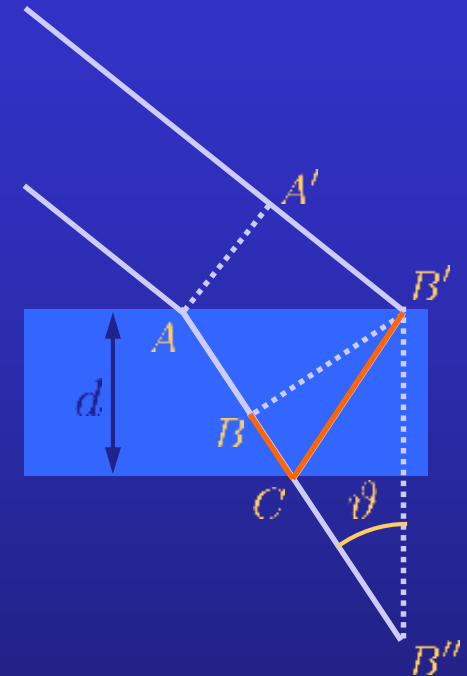
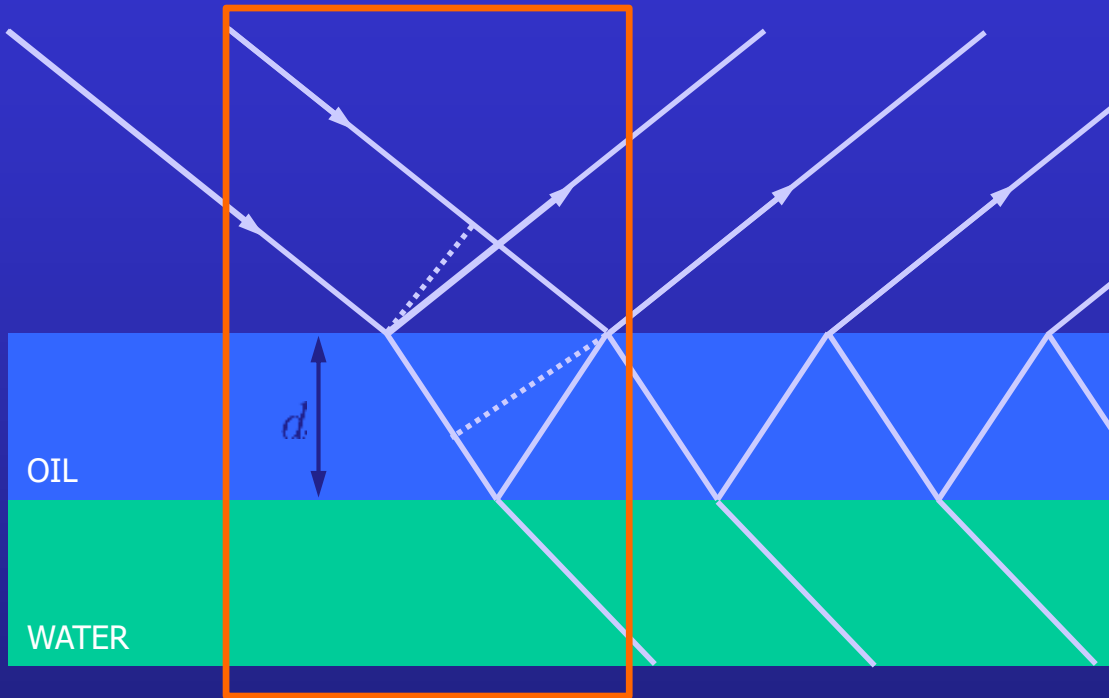


Thomas Poersch



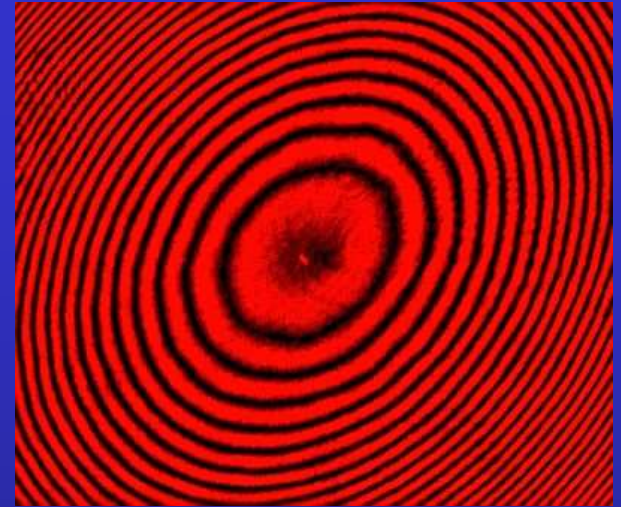
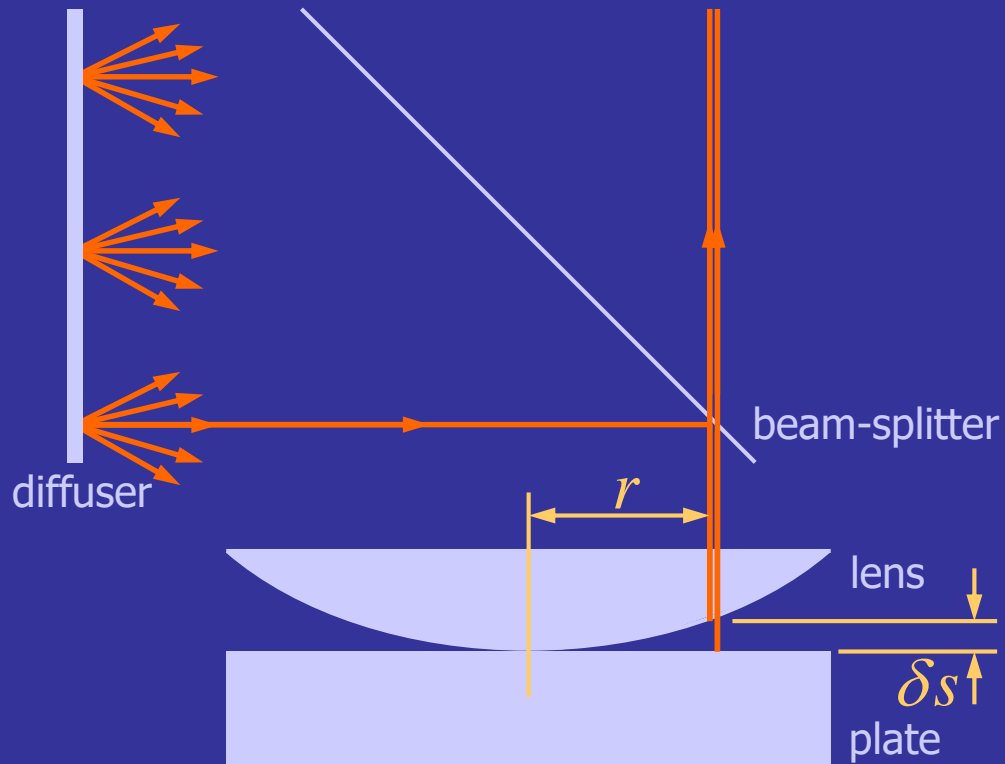
www.sciencebuddies.org

Thin film interference



- propagation time $T_{AB} = T_{A'B'}$
- path difference $BCB' = BCB'' = 2d \cos \vartheta$
- for constructive interference $2d \cos \vartheta = n\lambda$
...or destructive interference if phase changes upon reflection

Newton's rings



- apply to **sinusoidal waves** by taking into account the **phase** with which components arrive
- combine by adding the **amplitudes**
- contributions may therefore interfere **constructively** or **destructively**

Newton's rings

Interference occurs between the paths SAS and SBS, which differ in length by distance $2x$.

Pythagoras' theorem gives

$$\begin{aligned}R^2 &= r^2 + (R-x)^2 \\ &= r^2 + R^2 - 2Rx + x^2\end{aligned}$$

$$\Rightarrow r^2 = (2R-x)x$$

$$\approx 2Rx \quad \text{for } x \ll R$$

The phase difference between the two paths will be

$$\delta\phi = 2\pi \frac{2x}{\lambda} + \Delta\phi$$

where $\Delta\phi \equiv \phi_1 - \phi_2$ is the difference in phase change upon reflection from the two surfaces.

For one glass-air and one air-glass interface, $\Delta\phi = \pi$.

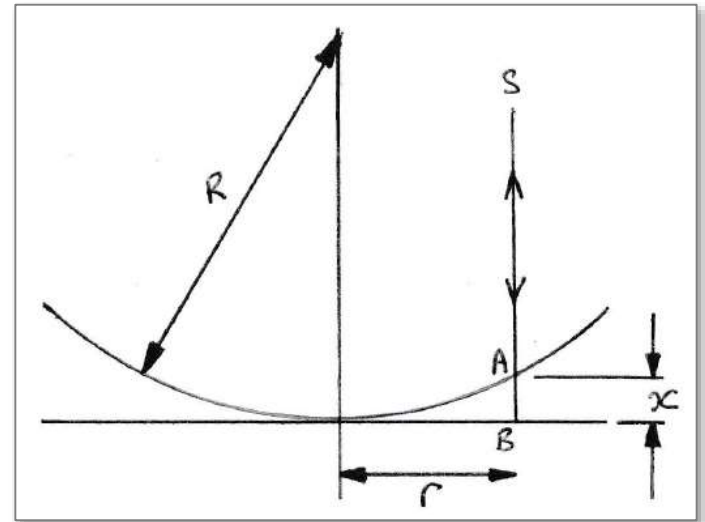
$$\Rightarrow \delta\phi = 4\pi \frac{x}{\lambda} + \pi$$

For a bright fringe, $\delta\phi = 2m\pi$; for a dark fringe, $\delta\phi = (2m+1)\pi$

$$\Rightarrow \text{dark fringes occur when } 4\pi \frac{x}{\lambda} + \pi = (2m+1)\pi$$

$$\therefore x = \frac{m}{2}\lambda$$

$$\Rightarrow r = \sqrt{2R \frac{m}{2}\lambda} = \sqrt{mR\lambda}$$



Wave Physics

PHYS 2023

Tim Freegarde

